

Review

In the previous lecture, we ...

- introduced the $\varepsilon - \delta$ identity
- introduced the concepts of metric tensors

Aims

In this lecture, we will ...

- introduce the conjugate metric tensor
- introduce the concepts of contravariant and covariant tensors

Example 2.11:

If the metric tensor g_{ij} is given by

$$g_{ij} = \frac{\partial z^k}{\partial x^i} \frac{\partial z^k}{\partial x^j},$$

then calculate the metric tensor for parabolic cylindrical coordinates (ξ, η, z) , where

$$x = \xi\eta, \quad y = \frac{1}{2}(\xi^2 - \eta^2), \quad z = z.$$



Answer:

$$g_{11} = \frac{\partial z^k}{\partial x^1} \frac{\partial z^k}{\partial x^1} = \left(\frac{\partial x}{\partial \xi} \right)^2 + \left(\frac{\partial y}{\partial \xi} \right)^2 + \left(\frac{\partial z}{\partial \xi} \right)^2 = \eta^2 + \xi^2.$$

$$g_{12} = \frac{\partial z^k}{\partial x^1} \frac{\partial z^k}{\partial x^2} = \frac{\partial x}{\partial \xi} \frac{\partial x}{\partial \eta} + \frac{\partial y}{\partial \xi} \frac{\partial y}{\partial \eta} + \frac{\partial z}{\partial \xi} \frac{\partial z}{\partial \eta} = \eta\xi - \xi\eta + 0 = 0.$$

and so forth. Finally, we obtain

$$g_{ij} = \begin{pmatrix} \eta^2 + \xi^2 & 0 & 0 \\ 0 & \eta^2 + \xi^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$



Exercise 2.12:

Consider the modified cylindrical polar coordinates (r, θ, z) where

$$x = r \cos \theta, \quad y = r \sin \theta, \quad z = z.$$

Find the metric tensor.



Answer:



Note:

In this exercise, not all of the off-diagonal elements are zero, i.e., $g_{12} = g_{21} = -r \sin \theta$. But even so, the metric tensor is *still* symmetric about the diagonal.

**2.3.3 Conjugate metric tensor**

For reasons seen a little bit later, sometimes it is more convenient to use the inverse of the metric tensor. This inverse is called the *conjugate metric tensor*, and is denoted by g^{jk} .

As g^{jk} is the inverse of g_{ij} then

$$g_{ij} g^{jk} = \delta_i^k.$$

To calculate g^{jk} , the usual method is to use *cofactors*, namely

$$g^{jk} = \frac{\text{cofactor}(g_{jk})}{g},$$

where g is the determinant of g_{ij} , i.e.,

$$g = \begin{vmatrix} \frac{\partial z^k}{\partial x^i} & \frac{\partial z^k}{\partial x^j} \\ \frac{\partial z^k}{\partial x^i} & \frac{\partial z^k}{\partial x^j} \end{vmatrix} = \begin{vmatrix} g_{11} & g_{12} & g_{13} \\ g_{21} & g_{22} & g_{23} \\ g_{31} & g_{32} & g_{33} \end{vmatrix}.$$

Question?

What is a cofactor?



Answer: (From MATH187)

**Example 2.12:**

Previously, for the modified cylindrical polar coordinates (r, θ, z) , where

$$x = r \cos \theta + r, \quad y = r \sin \theta, \quad z = z,$$

we found the metric tensor

$$g_{ij} = \begin{pmatrix} 2 \cos \theta + 2 & -r \sin \theta & 0 \\ -r \sin \theta & r^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

To find the conjugate metric tensor g^{jk} , we first need the determinant g , namely

$$g = \begin{vmatrix} 2 \cos \theta + 2 & -r \sin \theta & 0 \\ -r \sin \theta & r^2 & 0 \\ 0 & 0 & 1 \end{vmatrix} = r^2(2 \cos \theta - \sin^2 \theta + 2) = r^2(\cos \theta + 1)^2,$$

so that from the formulae

$$g^{jk} = \frac{(-1)^{j+k} \times G_{jk}}{g},$$

we find

$$g^{11} = \frac{(-1)^{1+1} \begin{vmatrix} r^2 & 0 \\ 0 & 1 \end{vmatrix}}{r^2(\cos \theta + 1)^2} = \frac{1}{(\cos \theta + 1)^2},$$

$$g^{12} = \frac{(-1)^{1+2} \begin{vmatrix} -r \sin \theta & 0 \\ 0 & 1 \end{vmatrix}}{r^2(\cos \theta + 1)^2} = \frac{\sin \theta}{r(\cos \theta + 1)^2},$$

etc.

Finally, we obtain

$$g^{jk} = \begin{pmatrix} \frac{1}{(\cos \theta + 1)^2} & \frac{\sin \theta}{r(\cos \theta + 1)^2} & 0 \\ \frac{\sin \theta}{r(\cos \theta + 1)^2} & \frac{2}{r^2(\cos \theta + 1)} & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

□

Exercise 2.13:

In Exercise 2.10, we found the metric tensor for cylindrical polar coordinates is

$$g_{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

Find the conjugate metric tensor. ❖

Answer:

□

Question?

In Exercise 2.13, there is an easier way to calculate g^{jk} . What is it? ❖

Answer:

□

Note:

A problem would occur if a diagonal element of g_{ij} was zero. However, from the definition of g_{ij} , this would only happen if the coordinate transformation

$$g_{\underline{i}\underline{i}} = (z^1_{,\underline{i}})^2 + (z^2_{,\underline{i}})^2 + (z^3_{,\underline{i}})^2 = 0 \quad \Rightarrow \quad z^k = z^k(x^1, x^2, x^3) = 0.$$

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2.3.4 Relationship between g and J

From section 2.3, recall that

$$J = \left| \frac{\partial X^i}{\partial x^j} \right| = \begin{vmatrix} \frac{\partial X^1}{\partial x^1} & \frac{\partial X^1}{\partial x^2} & \frac{\partial X^1}{\partial x^3} \\ \frac{\partial X^2}{\partial x^1} & \frac{\partial X^2}{\partial x^2} & \frac{\partial X^2}{\partial x^3} \\ \frac{\partial X^3}{\partial x^1} & \frac{\partial X^3}{\partial x^2} & \frac{\partial X^3}{\partial x^3} \end{vmatrix},$$

and if A and B are matrices, then $|AB| = |A||B|$ and $|A^T| = |A|$.

Thus, if we assume

$$X^i = z^i(x^1, x^2, x^3),$$

then from $J^2 = |J_{ij}| \times |J_{ij}| = |J_{ij}^T| \times |J_{ij}|$, we find ...

$$\begin{aligned} J^2 &= \left| \begin{pmatrix} \frac{\partial z^1}{\partial x^1} & \frac{\partial z^2}{\partial x^1} & \frac{\partial z^3}{\partial x^1} \\ \frac{\partial z^1}{\partial x^2} & \frac{\partial z^2}{\partial x^2} & \frac{\partial z^3}{\partial x^2} \\ \frac{\partial z^1}{\partial x^3} & \frac{\partial z^2}{\partial x^3} & \frac{\partial z^3}{\partial x^3} \end{pmatrix} \begin{pmatrix} \frac{\partial z^1}{\partial x^1} & \frac{\partial z^1}{\partial x^2} & \frac{\partial z^1}{\partial x^3} \\ \frac{\partial z^2}{\partial x^1} & \frac{\partial z^2}{\partial x^2} & \frac{\partial z^2}{\partial x^3} \\ \frac{\partial z^3}{\partial x^1} & \frac{\partial z^3}{\partial x^2} & \frac{\partial z^3}{\partial x^3} \end{pmatrix} \right|, \\ &= \begin{vmatrix} \frac{\partial z^k}{\partial x^1} \frac{\partial z^k}{\partial x^1} & \frac{\partial z^k}{\partial x^1} \frac{\partial z^k}{\partial x^2} & \frac{\partial z^k}{\partial x^1} \frac{\partial z^k}{\partial x^3} \\ \frac{\partial z^k}{\partial x^2} \frac{\partial z^k}{\partial x^1} & \frac{\partial z^k}{\partial x^2} \frac{\partial z^k}{\partial x^2} & \frac{\partial z^k}{\partial x^2} \frac{\partial z^k}{\partial x^3} \\ \frac{\partial z^k}{\partial x^3} \frac{\partial z^k}{\partial x^1} & \frac{\partial z^k}{\partial x^3} \frac{\partial z^k}{\partial x^2} & \frac{\partial z^k}{\partial x^3} \frac{\partial z^k}{\partial x^3} \end{vmatrix}, \\ &= g. \end{aligned}$$

Thus: $g = J^2$ or $J = \sqrt{g}$.

Note:

If $J = 0$, it means that there is no unique inverse, so that the metric tensor g_{ij} will have a zero determinant, i.e., $g = 0$.

If J is known, then it is easy to find g , and vice-versa.



Example 2.13:

In Example 2.13, we considered the coordinate transformation

$$x = r \cos \theta + r, \quad y = r \sin \theta, \quad z = z,$$

where we found

$$g = \left| \frac{dz^k}{dx^i} \frac{dz^k}{dx^j} \right| = r^2(\cos \theta + 1)^2.$$

To check, find J , i.e.,

$$\begin{aligned} J &= \left| \frac{\partial z^i}{\partial x^j} \right| = \begin{vmatrix} \cos \theta + 1 & -r \sin \theta & 0 \\ \sin \theta & r \cos \theta & 0 \\ 0 & 0 & 1 \end{vmatrix} = r \cos \theta (\cos \theta + 1) + r \sin^2 \theta \\ &= r(\cos \theta + 1). \end{aligned}$$

Thus,

$$J^2 = [r(\cos \theta + 1)]^2 = r^2(\cos \theta + 1)^2 = g.$$



Note:

We will see later that we need J in various equations, where we have already worked out g_{ij} . In such situations, it is much easier to calculate J using $J = \sqrt{g}$, than it is using $J = \left| \frac{dz^i}{dx^j} \right|$.



2.4 Classifying tensors

We now look at how to further classify tensors. To do this, we introduce the following definitions:

Definition 2.3:

A function $\phi(x^1, x^2, x^3)$ is called a *scalar* if upon transformation of coordinates $x^k = x^k(X^1, X^2, X^3)$ it does not change its original value, i.e.,

$$\phi(x^1(\tilde{X}), x^2(\tilde{X}), x^3(\tilde{X})) \equiv \Phi(X^1, X^2, X^3) = \phi(x^1, x^2, x^3).$$

□

Example 2.14:

An example of a scalar is the temperature at a point, for it has the same value in all coordinate systems.

□

Definition 2.4:

Quantities $A^k(x)$ are called *contravariant components of a vector* (or simply a *contravariant vector*) if under coordinate transformation $x^k = x^k(X^1, X^2, X^3)$ they transform according to the rule

$$\bar{A}^i(X) = A^j(x) \frac{\partial X^i}{\partial x^j}.$$

Here, \bar{A}^i are the components of the same vector, but with reference to coordinates X^j .

□

Example 2.15:

An example of a contravariant vector is the differential, i.e., $A^i = dx^i$, since according to the chain rule we have

$$dX^i = \frac{\partial X^i}{\partial x^j} dx^j.$$

□

Definition 2.5:

Similarly, $A_k(x)$ are called *covariant components of a vector* (or simply a *covariant vector*) if under coordinate transformation $x^k = x^k(X^1, X^2, X^3)$ they transform according to the rule

$$\bar{A}_i(X) = A_j(x) \frac{\partial x^j}{\partial X^i}.$$

□

Example 2.16:

An example of a covariant vector is the partial derivative of a scalar, i.e., $A_i = \frac{\partial \phi}{\partial x^i}$, since

$$\frac{\partial \phi}{\partial X^j} = \frac{\partial \phi}{\partial x^i} \frac{\partial x^i}{\partial X^j}.$$

□

Exercise 2.14:

For the following tensors, identify whether scalar, covariant or contravariant and write the transform transformations:

1. A^i **Answer:**
2. A
3. A_i
4. g_{ij}
5. g^{jk}
6. δ_j^i

□

Summary

In this lecture, we ...

- introduced the conjugate metric tensor
- introduced the concepts of contravariant and covariant tensors

Coming up

In the next lecture, we ...

- consider covariant/contravariant tensors of order n
- examine how to express the metric tensor in terms of base vectors

Homework Exercise 2.5:

1. For the spherical polar coordinate system (r, θ, ϕ) , which is related to the Cartesian coordinate system by

$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi, \quad z = r \cos \theta,$$

find the conjugate metric tensor.

2. Give an example, with appropriate transformation rules, of the following tensors:
 - (a) A scalar.
 - (b) A contravariant tensor.
 - (c) A covariant tensor.
 - (d) A mixed tensor.